

An analytical dynamic model of heat transfer from the heating body to the heated room

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Abstract. On the base of mathematical description of thermal balance the dynamic model of the hot-water heating body (radiator) was designed. The radiator is mathematically described as a heat transfer system between heating water and warmed-up air layer. Similarly, the dynamic model of heat transfer through the wall from the heated space to the outdoor environment was design. Both models were interconnected into dynamic model of heat transfer from the heating body to the heated room and they will be implemented into simulation model of the heating system in Matlab/Simulink environment.

1 Hot-water heating body

An analytic identification method based on mathematical description of a heat exchanger was chosen for design of the hot-water heating body dynamic model [1].

A hot-water heating body can be described as a heat transfer system between heating water and warmed-up air layer [2]. Heating water circulates inside radiator and delivers a heat through surface layer. External side of radiator is surrounded by air layer, which is warmed-up and heated air naturally flows up due to difference of specific weight [3]. It was considered ideal mixing of the heating water in internal space of the radiator and ideal air mixing in boundary layer of the radiator.

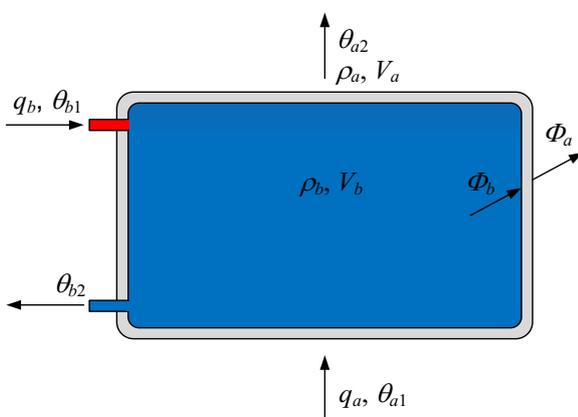


Fig. 1. Heat transfer in the heating body.

The heating water with mass flow q_b and temperature θ_{b1} inputs into radiator (Fig. 1). At every time point there is the volume of heating water V_b with density ρ_b in internal space of radiator. This water delivers a heat through surface layer of the radiator with plane A_r which is surrounded by air layer. Temperature of the return heating water from radiator is θ_{b2} . The air flow rate is q_a ,

volume of air in boundary layer of the radiator is V_a , its input temperature to the radiator surface is θ_{a1} and its output temperature after heating is θ_{a2} . Exchanged heat flow from heating water to the radiator wall is Φ_b and from the radiator wall to room space is Φ_a [4,5].

1.1 Mathematical description of heat transfer in the heating body

Generally, for dynamic balance describing of heat energy increasing or decreasing in system it is valid, that difference of input and output heat flows is equal to the time variation of the accumulated energy in a system [6]. Then the heat accumulation equation is valid for air heating in the boundary layer of the radiator:

$$\Phi_a - q_a \cdot c_a \cdot (\theta_{a2} - \theta_{a1}) = \rho_a \cdot V_a \cdot c_a \cdot \frac{d\theta_{a2}}{dt} \quad (1)$$

where Φ_a is heat flow from the radiator wall to the surrounding [W], q_a is air flow in the boundary layer of the radiator [$\text{kg} \cdot \text{s}^{-1}$], c_a is specific heat capacity (specific heat) of air [$\text{J} \cdot \text{kg}^{-1} \cdot \text{K}^{-1}$], θ_{a1} is input air temperature [K], θ_{a2} is output air temperature [K], ρ_a is air density [$\text{kg} \cdot \text{m}^{-3}$] and V_a is air volume in the boundary layer of the radiator [m^3].

For heat exchanging from the heating water to the radiator wall this heat accumulation equation is valid:

$$q_b \cdot c_b \cdot (\theta_{b1} - \theta_{b2}) - \Phi_b = \rho_b \cdot V_b \cdot c_b \cdot \frac{d\theta_{b2}}{dt} \quad (2)$$

where Φ_b is heat flow from the heating water to the radiator wall [W], q_b is mass flow rate of the heating water [$\text{kg} \cdot \text{s}^{-1}$], c_b is specific heat capacity (specific heat) of water [$\text{J} \cdot \text{kg}^{-1} \cdot \text{K}^{-1}$], θ_{b1} is input water temperature into the radiator [K], θ_{b2} is temperature of the return heating

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water from the radiator [K], ρ_b is water density [kg·m⁻³] and V_b is volume of water in the radiator [m³].

It is necessary for design of automatic control to know the dynamic dependence of the control deviation according to changes of variables which have effect on the deviation [7-10]. For that reason, we express dependent variables in equations (1) and (2) by their values in initial steady-state and their increments. Next, to simplify computation, we express each of dependences in non-dimensional form as relative changes of variables and we get system of differential equations describing the radiator:

$$x_{\phi a} - x_{\theta a 2} + x_{\theta a 1} - x_{q a} = \tau_{ma} \cdot \frac{dx_{\theta a 2}}{dt} \quad (3)$$

$$-x_{\theta b 2} + x_{\theta b 1} + x_{q b} - x_{\phi b} = \tau_{mb} \cdot \frac{dx_{\theta b 2}}{dt} \quad (4)$$

where $\tau_{ma} = \frac{\rho_a \cdot V_a}{q_{a0}}$ and $\tau_{mb} = \frac{\rho_b \cdot V_b}{q_{b0}}$ are time constants.

1.2 Dynamic model of the hot-water heating body

After Laplace transform of differential equations (3) and (4) and other mathematical operations it is valid for relative change of air output temperature:

$$X_{\theta a 2}(s) = \frac{1}{\tau_{ma} \cdot s + 1} [X_{\phi a}(s) + X_{\theta a 1}(s) - X_{q a}(s)] = G_a(s) \cdot [X_{\phi a}(s) + X_{\theta a 1}(s) - X_{q a}(s)] \quad (5)$$

$$X_{\theta b 2}(s) = \frac{1}{\tau_{mb} \cdot s + 1} [-X_{\phi b}(s) + X_{\theta b 1}(s) + X_{q b}(s)] = G_b(s) \cdot [-X_{\phi b}(s) + X_{\theta b 1}(s) + X_{q b}(s)] \quad (6)$$

where $\tau_{ma} = \frac{\rho_a \cdot V_a}{q_{a0}}$ and $\tau_{mb} = \frac{\rho_b \cdot V_b}{q_{b0}}$ are time constants.

To make mathematical description of dynamic characteristics of the radiator complete, it is necessary to supplement equation (5) and (6) with relations for relative changes of heat flows $x_{\phi a}$ and $x_{\phi b}$. If heat transfer coefficients h_a and h_b are constant in unsteady-states, relative variations of heat flows are dependent only on temperature changes of the heating water and air. Generally, we can express it as:

$$X_{\phi a}(s) = W_{ab}(s) \cdot X_{\theta b 2}(s) - W_{aa}(s) \cdot X_{\theta a 2}(s) \quad (7)$$

$$X_{\phi b}(s) = -W_{ba}(s) \cdot X_{\theta a 2}(s) + W_{bb}(s) \cdot X_{\theta b 2}(s) \quad (8)$$

Transfer functions in relations (7) and (8) are [2]:

$$W_{ab}(s) = \frac{k_b}{1 + \tau_r \cdot s} \quad W_{aa}(s) = k_a \cdot \frac{1 + \tau_b \cdot s}{1 + \tau_r \cdot s} \quad (9)$$

gains k_b, k_a :

$$k_b = \frac{A_r}{q_{b0} \cdot c_b} \cdot \frac{1}{\frac{1}{h_a} + \frac{1}{h_b} + \frac{d_r}{\lambda_r}} \quad (10)$$

$$k_a = \frac{A_r}{q_{a0} \cdot c_a} \cdot \frac{1}{\frac{1}{h_a} + \frac{1}{h_b} + \frac{d_r}{\lambda_r}} \quad (11)$$

and time constants τ_b, τ_r :

$$\tau_b = \rho_r \cdot c_r \cdot d_r \cdot \left(\frac{1}{h_b} + \frac{d_r}{2\lambda_r} \right) \quad (12)$$

$$\tau_r = \rho_r \cdot c_r \cdot d_r \cdot \frac{1 + \frac{(h_a + h_b) \cdot d_r}{2\lambda_r} + \frac{h_a \cdot h_b \cdot d_r}{6\lambda_r^2}}{h_a + h_b + \frac{h_a \cdot h_b \cdot d_r}{\lambda_r}} \quad (13)$$

where other variables not mentioned above are: A_r is plane of the radiator surface [m²], h_a is heat transfer coefficient between radiator surface and air [W·m⁻²·K⁻¹], h_b is heat transfer coefficient between heating water and radiator body [W·m⁻²·K⁻¹], d_r is thickness of the radiator wall [m], λ_r is heat conductivity coefficient of the radiator body material [W·m⁻¹·K⁻¹], ρ_r is specific weight of the radiator body material [kg·m⁻³] and c_r is specific heat capacity (specific capacity) of the radiator body material [J·kg⁻¹·K⁻¹] [11,12].

Transfer functions $W_{ba}(s)$ and $W_{bb}(s)$ we get by reciprocity of indexes „a“ and „b“ in terms (9).

Based on the terms (5) till (9) for non-dimensional variables the block diagram of dynamic model of the hot-water heating body was designed (Fig. 2).

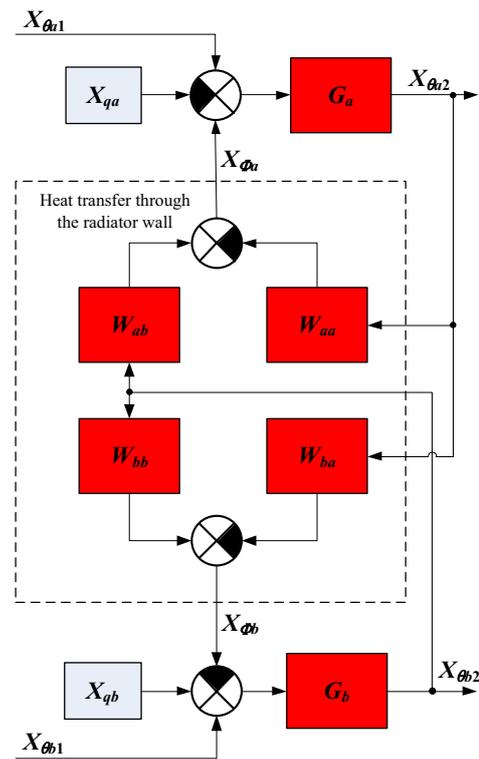


Fig. 2. Block diagram of dynamic model of the hot-water heating body.

Input variables into model are non-dimensional variables or constants: input temperature of air $X_{\theta a1}$, air flow rate X_{qa} , temperature of the input heating water into radiator $X_{\theta b1}$ and mass flow rate of the heating water X_{qb} . Output variables from model are non-dimensional variables: output temperature of air after heating by the radiator surface $X_{\theta a2}$ and temperature of the return heating water from the radiator $X_{\theta b2}$.

2 Heat transfer through the wall

For design of the dynamic model of heat transfer through the wall we have considered a plane wall, where the wall has been considered as continuum with continuously distributed thermal resistance and capacity [13,14]. We have chosen elementary layer with following parameters (Fig. 3): thickness of the plane wall d_w [m], layer thickness in the plane wall dy [m], distance of layer from the heated surface y [m]. The heat flow Φ_1 [W] is supplied into heated wall surface which temperature is θ_{w1} [K], the heat flow Φ_2 [W] is taken away from the refrigerated wall surface which temperature is θ_{w2} [K]. The heat flow Φ inputs into the unit surface of layer dy , and the heat flow $\Phi + d\Phi$ outputs from it. Temperature of the elementary layer is θ_w [K].

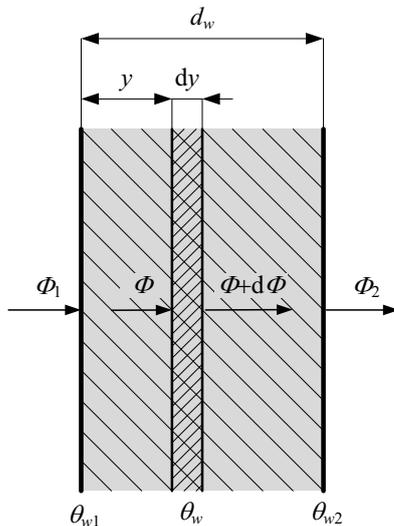


Fig. 3. Heat transfer through the plane wall.

2.1 Mathematical description of heat transfer through the wall

The heat energy does not originate either does not lose in considered elementary layer of the wall. Then difference of the input heat and output heat in the layer has to be equal to the time variation of the energy in the layer [15].

Let's c is specific heat capacity (specific heat) and ρ is volume weight of the wall material, then:

$$\Phi - (\Phi + d\Phi) = \frac{\partial}{\partial t} (c \cdot \rho \cdot \theta_w \cdot dy) \quad (14)$$

Considering that heat flow $d\Phi$ is:

$$d\Phi = \frac{\partial \Phi}{\partial y} dy \quad (15)$$

If specific heat capacity c and volume weight ρ of the used wall material are constant, then [16]:

$$-\frac{\partial \Phi}{\partial y} = c \cdot \rho \cdot \frac{\partial \theta_w}{\partial t} \quad (16)$$

According to Fourier's law the heat flow is directly proportional to the temperature gradient:

$$\Phi = -\lambda \frac{\partial \theta_w}{\partial y} \quad (17)$$

where λ is heat conductivity coefficient of the used wall material.

Partial differential equations (16) and (17) with relevant initial and border conditions completely describe non-stationary one-dimensional heat flow [17].

As we mentioned above we have expressed dependent variables by their values and increments and by substitutions and subtractions we have got partial differential equations system of heat transfer dynamics through the wall:

$$-\frac{\partial \Delta \Phi}{\partial y} = c \cdot \rho \cdot \frac{\partial \Delta \theta_w}{\partial t} \quad (18)$$

$$\Delta \Phi = -\lambda \frac{\partial \Delta \theta_w}{\partial y} \quad (19)$$

To simplify computation, we have expressed each of dependences in non-dimensional form and then we can express partial differential equations (18) and (19) in a form:

$$\frac{\partial x_\Phi}{\partial y} + c \cdot \rho \cdot \frac{d_w}{\lambda} \frac{\partial x_{\theta_w}}{\partial t} = 0 \quad (20)$$

$$x_\Phi + d_w \frac{\partial x_{\theta_w}}{\partial y} = 0 \quad (21)$$

Equations (16), (17) or (20), (21) still need to be supplemented by equations of heat transfer on both sides of the wall surfaces. In dimensionless form for indoor surface area it is valid:

$$x_{\Phi 1} = x_{\Phi 1}^* - \kappa_1 \cdot x_{\theta_{w1}} = (x_{\theta_{a2}} - x_{\theta_{w1}}) \cdot h_1 - \kappa_1 \cdot x_{\theta_{w1}} \quad (22)$$

where $x_{\Phi 1}^*$ includes also external effects on the heat flow transfer into the wall (temperature changes or heat transfer coefficient changes, etc.), $\kappa_1 = \frac{h_1 \cdot d_w}{\lambda}$ and h_1 is heat transfer coefficient between air and wall surface [$\text{W} \cdot \text{m}^{-2} \cdot \text{K}^{-1}$].

Similarly, for outdoor surface area it is valid:

$$x_{\phi 2} = x_{\phi 2}^* + \kappa_2 \cdot x_{\theta w 2} = (x_{\theta w 2} - x_{\theta o}) \cdot h_2 + \kappa_2 \cdot x_{\theta w 2} \quad (23)$$

where $x_{\phi 2}^*$ includes external conditions of the heat flow transfer from the wall, $\kappa_2 = \frac{h_2 \cdot d_w}{\lambda}$ and h_2 is heat transfer coefficient between wall surface and air [$W \cdot m^{-2} \cdot K^{-1}$].

2.2 Dynamic model of heat transfer through the wall

By the Laplace transform of partial differential equations (20) and (21) and by the other mathematical operations it has been possible to get a system of equations, which describes dependence of non-dimensional variables for heat flows Φ_1 , Φ_2 and temperatures θ_{w1} , θ_{w2} .

Laplace transform image of partial differential equations system is:

$$\bar{X}_{\phi} = \frac{q}{q^2 - p} \bar{X}_{\phi 1} - \frac{p}{q^2 - p} \bar{X}_{\theta w 1} \quad (24)$$

$$\bar{X}_{\theta w} = -\frac{1}{q^2 - p} \bar{X}_{\phi 1} + \frac{q}{q^2 - p} \bar{X}_{\theta w 1} \quad (25)$$

Using inverse Laplace transform we have got following equations:

$$\bar{X}_{\phi} = \cosh\sqrt{p} \cdot \bar{X}_{\phi 1} - \sqrt{p} \cdot \sinh\sqrt{p} \cdot \bar{X}_{\theta w 1} \quad (26)$$

$$\bar{X}_{\theta w} = -\frac{\sinh\sqrt{p}}{\sqrt{p}} \cdot \bar{X}_{\phi 1} + \cosh\sqrt{p} \cdot \bar{X}_{\theta w 1} \quad (27)$$

Let denote $\bar{X}_{\phi} = \bar{X}_{\phi 2}$ and $\bar{X}_{\theta w} = \bar{X}_{\theta w 2}$. And simultaneously because that the unknowns variables are the heat flow Φ_2 taken away from the refrigerated wall surface and temperature θ_{w1} of the heated wall surface, we express $\bar{X}_{\phi 2}$ and $\bar{X}_{\theta w 1}$ [2]:

$$\bar{X}_{\phi 2} = \frac{1}{\cosh\sqrt{p}} \bar{X}_{\phi 1} - \sqrt{p} \cdot \operatorname{tgh}\sqrt{p} \cdot \bar{X}_{\theta w 2} \quad (28)$$

$$\bar{X}_{\theta w 1} = -\frac{\operatorname{tgh}\sqrt{p}}{\sqrt{p}} \bar{X}_{\phi 1} + \frac{1}{\cosh\sqrt{p}} \bar{X}_{\theta w 2} \quad (29)$$

with transfer functions:

$$G_1(p) = \frac{1}{\cosh\sqrt{p}} \quad (30)$$

$$G_2(p) = \sqrt{p} \cdot \operatorname{tgh}\sqrt{p} \quad (31)$$

$$G_3(p) = \frac{\operatorname{tgh}\sqrt{p}}{\sqrt{p}} \quad (32)$$

Next problem is to find function which corresponds to (26) and (27) using inverse Laplace transform. Due to the fact, that transfer functions (28)-

(30) are not defined as images, we have to use Taylor series and substitution $p = T_s \cdot s$. Thus we get transfer functions:

$$G_1(p) = \frac{24}{24 + 12T_s s + T_s s^2} \quad (33)$$

$$G_2(p) = \frac{4(6T_s s + T_s s^2)}{24 + 12T_s s + T_s s^2} \quad (34)$$

$$G_3(p) = \frac{120 + 20T_s s + T_s s^2}{5(24 + 12T_s s + T_s s^2)} \quad (35)$$

where $T_s = \frac{d_w^2 \cdot \rho \cdot c}{\lambda}$ is constant which is depended on the wall properties.

Based on the terms (28) till (35) for non-dimensional variables the block diagram of dynamic model of heat transfer through the wall was designed (Fig. 4).

Input variables into model are non-dimensional variables: temperature of the heated air in the room $X_{\theta w 2}$ and outdoor temperature $X_{\theta o}$. Output variables from model are non-dimensional variables of the wall temperatures $X_{\theta w 1}$ and $X_{\theta w 2}$ [18-20].

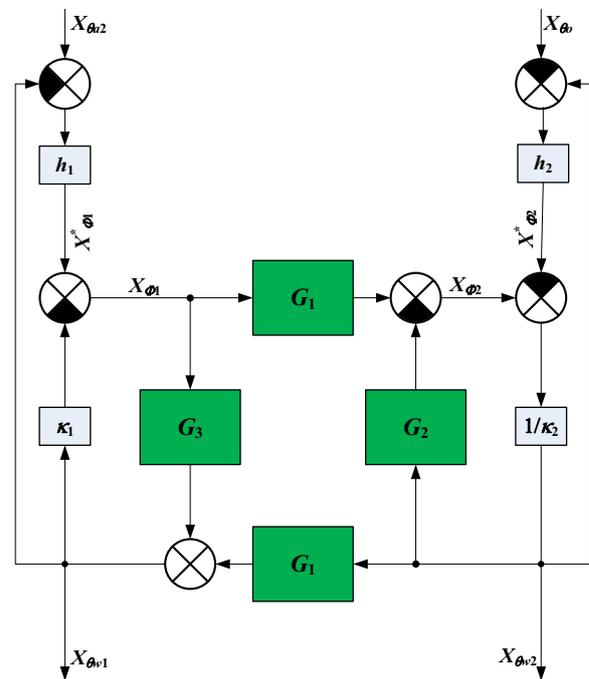


Fig. 4. Block diagram of dynamic model of heat transfer through the wall.

3 Conclusion

Dynamic model of heat transfer from the heating body to the heated room consists of two relative separated models: model of the heating body and model of heat transfer through the wall to the outdoor environment. Interconnection of these two models is presented in Fig. 5.

The designed dynamic model will be the basis for creating the simulation model of the heating system in Matlab/Simulink environment [21-25].

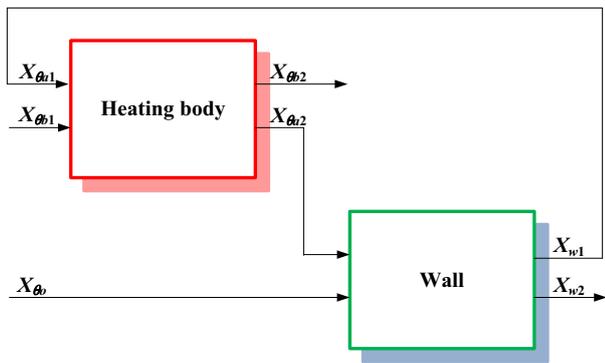


Fig. 5. Block diagram of dynamic model of heat transfer from the heating body to the heated room.

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