

Operational and Variational Formulations of Boundary Problems of Anisotropic Plate Analysis, Adapted for Application Within Discrete-Continual Methods

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Abstract. This paper is devoted to operational and variational formulations [1] of boundary problems of anisotropic plate analysis (with the use of so-called method of extended domain [2]), adapted for applications within discrete-continual methods (discrete-continual finite element method (DCFEM), discrete-continual variation-difference method (DCVDM)) [1]. Generally, the field of application of these methods, which are now becoming available for computer realization, comprises structures with regular (in particular, constant or piecewise constant) physical and geometrical parameters in some dimension. Considering problems remain continual along “basic” direction while along other directions discrete-continual methods presuppose mesh approximation.

1 Introduction

As is known, structures from anisotropic materials are widely used in modern construction. Anisotropic materials normally have sharp difference in the elastic properties for different directions. A classic vital sample here is natural wood. Its elastic modulus in tension parallel to the grain is well above corresponding modulus in tension across the grain. Besides, elastic constants of natural wood depend on the direction with respect to the wood grains. Anisotropic (and inhomogeneous) materials also include composite materials (particularly rather popular in aircraft engineering), crystals, some rocks, concrete, etc. Here we indicated materials with “natural” anisotropy. However, materials with so-called “artificial” anisotropy are applied in construction as well. We should mention here as a samples corrugated plates and shells from isotropic material or rib-

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reinforced plates and shell. Thus we have to consider boundary problems of anisotropic plate analysis [3-6].

2 Standard operational and variational formulations of boundary problem within method of extended domain

Operational formulation of the problem in the domain ω has the form:

$$\sum_{i=1}^2 \sum_{j=1}^2 \partial_j^2 \bar{B}_{i,j} \partial_i^2 w + 4\partial_1 \partial_2 \bar{B}_{3,3} \partial_1 \partial_2 w + 2 \sum_{i=1}^2 (\partial_i^2 \bar{B}_{i,3} \partial_1 \partial_2 w_k + \partial_1 \partial_2 \bar{B}_{i,3} \partial_i^2 w_k) + \bar{c}w = \theta q + \delta'_r M_r + \delta_r Q_r, \quad (1)$$

where Ω is the domain, occupied by structure with boundary $\Gamma = \partial\Omega$; ω is extended domain, embordering Ω ; x_1, x_2 are coordinates in plan; x_3 is coordinate of plate thickness; $\theta = \theta(x_1, x_2)$ is the characteristic function of domain Ω ; $\delta_r = \delta_r(x_1, x_2)$ is the delta-function of border $\Gamma = \partial\Omega$; $\bar{n} = [n_1 \ n_2]^T$ is unit normal vector of domain boundary $\Gamma = \partial\Omega$; $\delta'_r = \delta'_r(x_1, x_2)$ is corresponding normal derivative; w is plate deflection in domain Ω ; q is the load density in domain Ω ; Q_r, M_r are shear force and torque moment at the border $\Gamma = \partial\Omega$; c is modulus of foundation;

$$\partial_k = \partial / \partial x_k, \quad \partial_k^* = -\partial / \partial x_k, \quad k = 1, 2; \quad (2)$$

$$\bar{c} = \theta c; \quad \delta'_r = \delta'_r(x_1, x_2) = \partial \delta_r(x_1, x_2) / \partial \bar{n}; \quad (3)$$

$B_{1,1}, B_{2,2}, B_{3,3}, B_{1,2} = B_{2,1}, B_{2,3} = B_{3,2}, B_{1,3} = B_{3,1}$ are constants (parameters) that characterize the elastic properties of structure;

$$\bar{B}_{1,1} = \theta B_{1,1}, \quad \bar{B}_{1,3} = \theta B_{1,3}, \quad \bar{B}_{2,3} = \theta B_{2,3}, \quad \bar{B}_{1,2} = \theta B_{1,2}, \quad \bar{B}_{3,3} = \theta B_{3,3}; \\ \bar{B}_{1,2} = \theta B_{1,2}, \quad \bar{B}_{3,3} = \theta B_{3,3}, \quad \bar{B}_{2,2} = \theta B_{2,2}; \quad (4)$$

and in the particular case of an isotropic plate, we have

$$B_{1,1} = B_{2,2} = D; \quad B_{1,2} = B_{2,1} = D\nu; \quad B_{3,3} = 0.5 \cdot D(1 - \nu); \quad B_{1,3} = B_{2,3} = B_{3,1} = B_{3,2} = 0. \quad (5)$$

Internal moments and plate forces are determined by formulas

$$M_1 = B_{1,1}\chi_1 + B_{1,2}\chi_2 + B_{1,3}\chi_{12}; \quad M_2 = B_{1,2}\chi_1 + B_{2,2}\chi_2 + B_{2,3}\chi_{12}; \quad (6)$$

$$M_{12} = M_{21} = B_{1,3}\chi_1 + B_{2,3}\chi_2 + B_{3,3}\chi_{12}; \quad (7)$$

$$N_1 = B_{1,1}\partial_1\chi_1 + 3B_{1,3}\partial_2\chi_1 + (B_{1,2} + 2B_{3,3})\partial_1\chi_2 + B_{2,3}\partial_2\chi_2; \quad (8)$$

$$N_2 = B_{1,3}\partial_1\chi_1 + (B_{1,2} + 2B_{3,3})\partial_2\chi_1 + 3B_{2,3}\partial_1\chi_2 + B_{2,2}\partial_2\chi_2; \quad (9)$$

where $\chi_1, \chi_2, \chi_{12} = \chi_{21}$ are strains for $x_3 = 1$,

$$\chi_1 = -\partial_1^2 w; \quad \chi_2 = -\partial_2^2 w; \quad \chi_{12} = \chi_{21} = -2\partial_1 \partial_2 w \quad (10)$$

Variational formulation of the problem is formulated in the form of the corresponding energy functional (this formulation is convenient in the algorithmic sense):

$$\Phi(w) = \frac{1}{2} \iint_{\omega} [\theta(M_1 \chi_1 + M_2 \chi_2 + 2M_{1,2} \chi_{1,2}) + \bar{c} w^2] dx_1 dx_2 - \iint_{\omega} q w dx_1 dx_2. \quad (11)$$

3 Operational and variational displacement formulations of boundary problem within method of extended domain

If we combine (6)-(10) we can get formulation (1) in the following form

$$Lw = F, \quad (12)$$

where L is so-called stiffness operator,

$$L = \partial_1^2 \bar{B}_{1,1} \partial_1^2 + \partial_2^2 \bar{B}_{1,2} \partial_1^2 + \partial_1^2 \bar{B}_{2,1} \partial_2^2 + \partial_2^2 \bar{B}_{2,2} \partial_2^2 + 4\partial_1 \partial_2 \bar{B}_{3,3} \partial_1 \partial_2 + 2\partial_1^2 \bar{B}_{1,3} \partial_1 \partial_2 + 2\partial_2^2 \bar{B}_{2,3} \partial_1 \partial_2 + 2\partial_1 \partial_2 \bar{B}_{1,3} \partial_1^2 + 2\partial_1 \partial_2 \bar{B}_{2,3} \partial_2^2 + \bar{c}; \quad (13)$$

$$F = \theta q + \delta'_r M_r + \delta_r Q. \quad (14)$$

We can rewrite (11) in matrix form:

$$\Phi(w) = \frac{1}{2} \iint_{\omega} [(\theta \bar{N}, \bar{\chi}) + \bar{c} w^2] dx_1 dx_2 - \iint_{\omega} q w dx_1 dx_2, \quad (15)$$

where

$$\bar{N} = [M_1 \quad M_2 \quad M_{1,2} \quad M_{2,1}]^T; \quad \bar{\chi} = [\chi_1 \quad \chi_2 \quad \chi_{1,2} \quad \chi_{2,1}]^T. \quad (16)$$

Using (6)-(9), we get

$$\bar{N} = A \bar{\chi}, \quad \bar{\chi} = B w, \quad (17)$$

where

$$A = \begin{bmatrix} B_{1,1} & B_{1,2} & 0.5 \cdot B_{1,3} & 0.5 \cdot B_{1,3} \\ B_{1,2} & B_{2,2} & 0.5 \cdot B_{2,3} & 0.5 \cdot B_{2,3} \\ B_{1,3} & B_{2,3} & 0.5 \cdot B_{3,3} & 0.5 \cdot B_{3,3} \\ B_{1,3} & B_{2,3} & 0.5 \cdot B_{3,3} & 0.5 \cdot B_{3,3} \end{bmatrix}; \quad B = - \begin{bmatrix} \partial_1^2 \\ \partial_2^2 \\ 2\partial_1 \partial_2 \\ 2\partial_1 \partial_2 \end{bmatrix}; \quad (18)$$

Thus we have:

$$\Phi(w) = \frac{1}{2} \iint_{\omega} [(B^* \bar{A} B w, w) + \bar{c} w^2] dx_1 dx_2 - \iint_{\omega} q w dx_1 dx_2, \quad (19)$$

where

$$\bar{A} = \theta A = \begin{bmatrix} \bar{B}_{1,1} & \bar{B}_{1,2} & 0.5 \cdot \bar{B}_{1,3} & 0.5 \cdot \bar{B}_{1,3} \\ \bar{B}_{1,2} & \bar{B}_{2,2} & 0.5 \cdot \bar{B}_{2,3} & 0.5 \cdot \bar{B}_{2,3} \\ \bar{B}_{1,3} & \bar{B}_{2,3} & 0.5 \cdot \bar{B}_{3,3} & 0.5 \cdot \bar{B}_{3,3} \\ \bar{B}_{1,3} & \bar{B}_{2,3} & 0.5 \cdot \bar{B}_{3,3} & 0.5 \cdot \bar{B}_{3,3} \end{bmatrix}; \quad B^* = -[\partial_1^2 \quad \partial_2^2 \quad 2\partial_2 \partial_1 \quad 2\partial_2 \partial_1]. \quad (20)$$

4 Operational and variational formulations of boundary problem in displacements with extraction of basic direction

Without loss of generality we suppose constancy of physical and geometrical parameters of plate along coordinate x_2 (it is so-called “basic direction”).

We can rewrite (12) in the following form:

$$L = -L_4 \partial_2^4 + L_3 \partial_2^3 + L_2 \partial_2^2 + L_1 \partial_2 + L_0 + \bar{c}, \quad (21)$$

where

$$L_4 = -\bar{B}_{2,2}; \quad L_3 = 2\bar{B}_{2,3} \partial_1 + 2\partial_1 \bar{B}_{2,3}; \quad L_2 = \bar{B}_{1,2} \partial_1^2 + \partial_1^2 \bar{B}_{2,1} + 4\partial_1 \bar{B}_{3,3} \partial_1; \\ L_1 = 2\partial_1^2 \bar{B}_{1,3} \partial_1 + 2\partial_1 \bar{B}_{1,3} \partial_1^2; \quad L_0 = \partial_1^2 \bar{B}_{1,1} \partial_1^2. \quad (22)$$

We can rewrite (21):

$$-L_4 \partial_2^4 w + L_3 \partial_2^3 w + L_2 \partial_2^2 w + L_1 \partial_2 w + (L_0 + \bar{c})w = F. \quad (23)$$

Let us introduce the following notation

$$y_1 = y_1(x_1, x_2) = w(x_1, x_2); \quad y_i = y_i(x_1, x_2) = \partial_2^{i-1} w(x_1, x_2), \quad i = 2, 3, 4. \quad (24)$$

If we combine (23) and (24) we get:

$$-L_4 \partial_2 y_4 + L_3 y_4 + L_2 y_3 + L_1 y_2 + (L_0 + \bar{c})y_1 = F. \quad (25)$$

Uniting (24) and (25) we obtain

$$\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & L_4 \end{bmatrix} \begin{bmatrix} y_1' \\ y_2' \\ y_3' \\ y_4' \end{bmatrix} = \begin{bmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ L_0 + \bar{c} & L_1 & L_2 & L_3 \end{bmatrix} \begin{bmatrix} y_1 \\ y_2 \\ y_3 \\ y_4 \end{bmatrix} - \begin{bmatrix} 0 \\ 0 \\ 0 \\ F \end{bmatrix}, \quad (26)$$

where

$$y_i'(x_1, x_2) = \partial_2 y_i(x_1, x_2), \quad i = 2, 3, 4. \quad (27)$$

Thus we have the following system of the first-order ordinary differential (with respect to x_2) equations with operator coefficients

$$\begin{bmatrix} y_1' \\ y_2' \\ y_3' \\ y_4' \end{bmatrix} = \begin{bmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ L_4^{-1}(L_0 + \bar{c}) & L_4^{-1}L_1 & L_4^{-1}L_2 & L_4^{-1}L_3 \end{bmatrix} \begin{bmatrix} y_1 \\ y_2 \\ y_3 \\ y_4 \end{bmatrix} - \begin{bmatrix} 0 \\ 0 \\ 0 \\ L_4^{-1}F \end{bmatrix} \quad (28)$$

or

$$\bar{U}' = \tilde{L}\bar{U} + \tilde{F}, \quad (29)$$

where

$$\tilde{L} = \begin{bmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ L_4^{-1}(L_0 + \bar{c}) & L_4^{-1}L_1 & L_4^{-1}L_2 & L_4^{-1}L_3 \end{bmatrix}; \quad \tilde{F} = - \begin{bmatrix} 0 \\ 0 \\ 0 \\ L_4^{-1}F \end{bmatrix}; \quad (30)$$

$$\bar{U} = \begin{bmatrix} y_1 \\ y_2 \\ y_3 \\ y_4 \end{bmatrix}; \quad \bar{U}' = \partial_2 \bar{U} = \begin{bmatrix} \partial_2 y_1 \\ \partial_2 y_2 \\ \partial_2 y_3 \\ \partial_2 y_4 \end{bmatrix} = \begin{bmatrix} y'_1 \\ y'_2 \\ y'_3 \\ y'_4 \end{bmatrix}. \quad (31)$$

Equations (29), of course, should be supplemented with boundary conditions given in cross-sections with coordinates $x_{2,k}^b$, $k = 1, \dots, n_k$. This boundary conditions have form

$$B_k^- \bar{U}_{k-1}(x_{2,k}^b - 0) + B_k^+ \bar{U}_k(x_{2,k}^b + 0) = \bar{g}_k^- + \bar{g}_k^+, \quad k = 2, \dots, n_k - 1; \quad (32)$$

$$B_1^+ \bar{U}_1(x_{2,1}^b + 0) + B_{n_k}^- \bar{U}_{n_k-1}(x_{2,n_k}^b - 0) = \bar{g}_1^+ + \bar{g}_{n_k}^-, \quad (33)$$

where B_k^-, B_k^+ , $k = 2, \dots, n_k - 1$, $B_1^+, B_{n_k}^-$ are matrices of boundary conditions of the fourth order; \bar{g}_k^-, \bar{g}_k^+ , $k = 2, \dots, n_k - 1$, $\bar{g}_1^+, \bar{g}_{n_k}^-$ are right-side vectors of boundary conditions.

Operators (18) and (21) can be rewritten in the following form:

$$B = \partial_1^2 B_{2,0} + \partial_1 \partial_2 B_{1,1} + \partial_2^2 B_{0,2}; \quad B^* = \partial_1^2 B_{2,0}^* + \partial_1 \partial_2 B_{1,1}^* + \partial_2^2 B_{0,2}^*, \quad (34)$$

where

$$B_{2,0} = -[1 \ 0 \ 0 \ 0]^T; \quad B_{1,1} = -[0 \ 0 \ 2 \ 2]^T; \quad B_{0,2} = -[0 \ 1 \ 0 \ 0]^T; \quad (35)$$

$$B_{2,0}^* = -[1 \ 0 \ 0 \ 0]; \quad B_{1,1}^* = -[0 \ 0 \ 2 \ 2]; \quad B_{0,2}^* = -[0 \ 1 \ 0 \ 0]. \quad (36)$$

Thus, we have:

$$\begin{aligned} (B^* \bar{A} B w, w) &= ((\partial_1^2 B_{2,0}^* + \partial_2^2 \partial_1 B_{1,1}^* + \partial_2^2 B_{0,2}^*) \bar{A} (\partial_1^2 B_{2,0} + \partial_1 \partial_2 B_{1,1} + \partial_2^2 B_{0,2}) w, w) = \\ &= (\partial_1^2 B_{2,0}^* \bar{A} \partial_1^2 B_{2,0} w, w) + (\partial_1^2 B_{2,0}^* \bar{A} \partial_1 \partial_2 B_{1,1} w, w) + \\ &+ (\partial_1^2 B_{2,0}^* \bar{A} \partial_2^2 B_{0,2} w, w) + (\partial_2^2 \partial_1 B_{1,1}^* \bar{A} \partial_1^2 B_{2,0} w, w) + \\ &+ (\partial_2^2 \partial_1 B_{1,1}^* \bar{A} \partial_1 \partial_2 B_{1,1} w, w) + (\partial_2^2 \partial_1 B_{1,1}^* \bar{A} \partial_2^2 B_{0,2} w, w) + \\ &+ (\partial_2^2 B_{0,2}^* \bar{A} \partial_1^2 B_{2,0} w, w) + (\partial_2^2 B_{0,2}^* \bar{A} \partial_1 \partial_2 B_{1,1} w, w) + \\ &+ (\partial_2^2 B_{0,2}^* \bar{A} \partial_2^2 B_{0,2} w, w) = (B_{0,2}^* \bar{A} B_{0,2} y_3, y_3) + \\ &+ (B_{0,2}^* \bar{A} \partial_1 B_{1,1} y_2, y_3) + (\partial_1 B_{1,1}^* \bar{A} B_{0,2} y_3, y_2) + \\ &+ (B_{0,2}^* \bar{A} B_{2,0} \partial_1^2 y_1, y_3) + (\partial_1^2 B_{2,0}^* \bar{A} B_{0,2} y_3, y_1) + \\ &+ (\partial_1 B_{1,1}^* \bar{A} B_{2,0} \partial_1^2 y_1, y_2) + (\partial_1^2 B_{2,0}^* \bar{A} \partial_1 B_{1,1} y_2, y_1) + \\ &+ (\partial_1 B_{1,1}^* \bar{A} B_{1,1} \partial_1 y_2, y_2) + (\partial_1^2 B_{2,0}^* \bar{A} B_{2,0} \partial_1^2 y_1, y_1), \end{aligned}$$

and finally

$$\begin{aligned} (B^* \bar{A} B w, w) &= (B_{0,2}^* \bar{A} B_{0,2} y_3, y_3) + (B_{0,2}^* \bar{A} \tilde{B}_{1,1} y_2, y_3) + (\tilde{B}_{1,1}^* \bar{A}_k B_{0,2} y_3, y_2) + \\ &+ (B_{0,2}^* \bar{A} \tilde{B}_{2,0} y_1, y_3) + (\tilde{B}_{2,0}^* \bar{A} B_{0,2} y_3, y_1) + (\tilde{B}_{1,1}^* \bar{A} \tilde{B}_{2,0} y_1, y_2) + \\ &+ (\tilde{B}_{2,0}^* \bar{A} \tilde{B}_{1,1} y_2, y_1) + (\tilde{B}_{1,1}^* \bar{A} \tilde{B}_{1,1} y_2, y_2) + (\tilde{B}_{2,0}^* \bar{A} \tilde{B}_{2,0} y_1, y_1), \end{aligned} \quad (37)$$

where

$$B_{2,0} = -[\partial_1^2 \quad 0 \quad 0 \quad 0]^T; \quad B_{1,1} = -[0 \quad 0 \quad 2\partial_1 \quad 2\partial_1]^T; \quad (38)$$

$$B_{2,0}^* = -[\partial_1^2 \quad 0 \quad 0 \quad 0]; \quad B_{1,1}^* = -[0 \quad 0 \quad 2\partial_1^* \quad 2\partial_1^*]. \quad (39)$$

We have

$$(B^* \bar{A} B w, w) = (K \bar{U}^r, \bar{U}^r), \quad (40)$$

where

$$K = \begin{bmatrix} \tilde{B}_{2,0}^* \bar{A} \tilde{B}_{2,0} & \tilde{B}_{2,0}^* \bar{A} \tilde{B}_{1,1} & \tilde{B}_{2,0}^* \bar{A} B_{0,2} \\ \tilde{B}_{1,1}^* \bar{A} \tilde{B}_{2,0} & \tilde{B}_{1,1}^* \bar{A} \tilde{B}_{1,1} & \tilde{B}_{1,1}^* \bar{A}_k B_{0,2} \\ B_{0,2}^* \bar{A} \tilde{B}_{2,0} & B_{0,2}^* \bar{A} \tilde{B}_{1,1} & B_{0,2}^* \bar{A} B_{0,2} \end{bmatrix}; \quad \bar{U}^r = \begin{bmatrix} y_1 \\ y_2 \\ y_3 \end{bmatrix}. \quad (41)$$

Thus, functional (19) can be rewritten in the following form:

$$\Phi(w) = \Phi(\bar{U}^r) = \frac{1}{2} \iint_{\omega} (\tilde{K} \bar{U}^r, \bar{U}^r) dx_1 dx_2 - \iint_{\omega} (\bar{F}, \bar{U}^r) w dx_1 dx_2, \quad (42)$$

where

$$\tilde{K} = \begin{bmatrix} \tilde{B}_{2,0}^* \bar{A} \tilde{B}_{2,0} + \bar{c} & \tilde{B}_{2,0}^* \bar{A} \tilde{B}_{1,1} & \tilde{B}_{2,0}^* \bar{A} B_{0,2} \\ \tilde{B}_{1,1}^* \bar{A} \tilde{B}_{2,0} & \tilde{B}_{1,1}^* \bar{A} \tilde{B}_{1,1} & \tilde{B}_{1,1}^* \bar{A}_k B_{0,2} \\ B_{0,2}^* \bar{A} \tilde{B}_{2,0} & B_{0,2}^* \bar{A} \tilde{B}_{1,1} & B_{0,2}^* \bar{A} B_{0,2} \end{bmatrix}; \quad \bar{F} = \begin{bmatrix} F \\ 0 \\ 0 \end{bmatrix}. \quad (43)$$

The solution of this problem is the point (function) of the constrained extremum of this functional with the condition (27), taking into account (32), (33).

These formulations are the initial ones for the realization of the discrete-continual methods for the problems of anisotropic plate analysis.

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