

Comments on Variational Method and Energy Method in Computational Mechanics

Xingbo WANG Biao LU and Ying TIAN

Department of Mechatronics, Foshan University, Foshan City, PRC, 528000[1-3]

Abstract. The article demonstrates that the energy method utilized in computational mechanics is a special form of the variational method through mathematical deduction and comparison to the two methods. The result discloses a basic fact that mechanical computations performed by different methods are intrinsically related in some way. The article also points out that a deep study on this intrinsic relation is necessary and helpful to guarantee loyalty of the computational results to the computed objectives.

1 Introduction

Computational mechanics is a subject that combines various skills and techniques in both mechanics and computer science. Ever since the being of the subject, people have developed tens of methods to perform kinds of mechanical computations, e.g., boundary element method (BEM)[1], the finite element method (FEM)[2], the finite volume method (FVM)[3], the collocation method (CM)[4], Numerical manifold method (NMM)[5], the meshless method or element free method (EFM)[6] and many other methods as are introduced in reference [7]. Literatures[8]-[10] also show that the Galerkin method(GM) , Ritz method (or Rayleigh–Ritz method) (RM) are frequently applied to derive all the BEM, FEM, NMM and so on. Also it is can see that, the energy methods (E_M)[11][12,13], the variational method (VariM)[14] and the weighted residual method (WRM) [15] are considered fundamental principles in computational mechanics. Since the ultimate purpose of computational mechanics is to solve numerically a partial differential equation (PDE) that is derived from a mechanical problem, I present in reference [16] a hierarchy diagram, as shown in Fig 1, to classify the classes of the methods describe the relationship of the methods mentioned above.

In the diagram, one can see that, either of the GM and the RM can deduces FEM, BEM, EFGM, or FDM while the GM and the RM can have different sources: the GM is deduced only from the WRM while he RM can be derived from every one of the WRM, the VariM and the E_M.

Now come at least three questions in front of the engineers:

(1). Are the GM and the RM identical ones?

(2). Are the WRM, the VariM and the E_M identical ones?

(3). If the GM and the RM are not identical or the WRM, the VariM and the E_M identical are not identical, which one is sure?

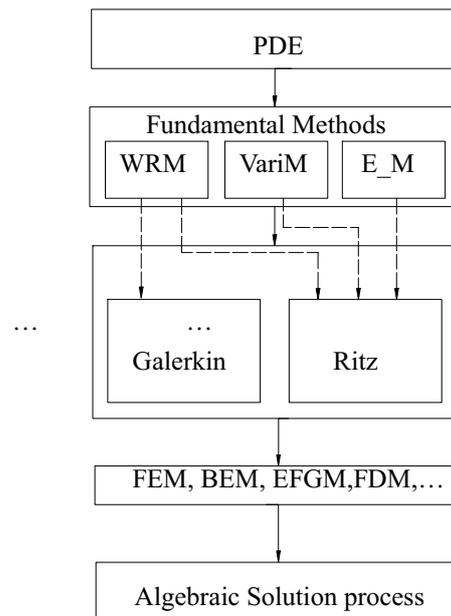


Fig 1. Hierarchy diagram of numerical methods to solve PDEs

Weile DS and Wildman RA might answer the question (1) in their article [8]. They say ” the Rayleigh–Ritz and Galerkin methods are in principle very different, they often lead to the same set of discretized equations, are closely related, and are often discussed as if they were identical.” But what are the answers to the reminders? Not finding an explicit answer, this paper proves theoretically that the E_M is identical to VariM under

some special conditions. That is to say, the E_M is a special VariM.

2 From Variational Method to Energy Method

The VariM are methods based on the variational principle. In order to show that the E_M is a special VariM, we begin with the general principle of variational method, and then finish the proofs.

2.1 Equivalent variational form of a PDE

Given a PDE

$$L\phi - f = 0 \quad \text{in } \Omega \quad (1)$$

where $L = -(\sum_{i,j=1}^2 \partial_j (a_{ij}(x_1, x_2) \partial_i) + a_{00}(x_1, x_2))$

is an operator of 2nd order PDE, and Ω is the domain of the function ϕ and f .

Multiplying (1) by a variation $\delta\phi$, and integrate over the domain Ω , it yields

$$\int_{\Omega} \underbrace{d\Omega(L\phi - f)}_{PDE} \delta\phi = 0 \quad (2)$$

Then integrate by part is used to transfer the derivatives in the operator L to $\delta\phi$, and the δ symbol is brought outside the integral and obtain an equation in the form

$$\delta \left(\int_{\Omega} \Phi(\phi) d\Omega \right) = 0 \quad (3)$$

Denote (3) by

$$\delta I(\phi) = 0 \quad (4)$$

and the equation (4) is the equivalent variational form of the PDE (1) in which $I(\phi)$ is the variational form of the PDE.

The equation (4) is a general form of general meaning. When applied in concrete problem, it may have definite form. For example, if L be a symmetric, positive definite, real operator, and a PDE is with Dirichlet boundary condition, as shown by

$$\begin{cases} L\phi = f & \text{in } \Omega \\ \phi = g & \text{on } \partial\Omega \end{cases} \quad (5)$$

the variational form of the PDE is

$$I(\phi) = \frac{1}{2} \langle L\phi, \phi \rangle - \langle \phi, f \rangle \quad (6)$$

If the PDE is with Neumann boundary condition

$$\begin{cases} L\phi = f & \text{in } \Omega \\ \sum_{i,j} a_{i,j} \partial_n \phi = h & \text{on } \partial\Omega \end{cases} \quad (7)$$

then its variational form is

$$I(\phi) = \frac{1}{2} \langle L\phi, \phi \rangle - \langle \phi, f \rangle - \langle h, \phi \rangle_{\partial\Omega} \quad (8)$$

where

$$\langle Lu, v \rangle = \int_{\Omega} (Lu)v d\Omega, \langle f, v \rangle = \int_{\Omega} f v d\Omega, \langle h, v \rangle_{\partial\Omega} = \int_{\partial\Omega} h v d(\partial\Omega)$$

The equations (6) and (9) are fundamental equations to deduce Rayleigh-Ritz Method, for which we omit details here.

2.2 Energy Methods

In mechanics, the energy methods refer to either the virtual work (or virtual displacement) method, or the complementary virtual work method. Now we prove that both methods are special forms of VariM.

2.2.1 Virtual Work Method (VWM)

Virtual work method comes from is also called principle of virtual work (PVW) or principle of virtual displacement (PVD). Virtual work δW is work done by a force F along a virtual displacement δd

$$\delta W = F \cdot \delta d \quad (9)$$

In a generalized coordinate system of N generalized coordinates (q_1, q_2, \dots, q_N) , the PVM says

$$\delta W = \sum_{i=1}^N Q_i \delta q_i = 0 \quad (10)$$

where $\{Q_i\}_{i=1}^N$ are the generalized forces acting on the system, $\{\delta q_i\}_{i=1}^N$ are virtual displacements and δW are the virtual work done by $\{Q_i\}_{i=1}^N$ along $\{\delta q_i\}_{i=1}^N$.

Now consider in equation (9), $F(r)$ is a conservative force along a path r , then work done by $F(r)$ from r_1 to r_2 on the path is

$$W = \int_{r_1}^{r_2} F(r) \cdot dr = W(r_2) - W(r_1) = W(r_1, r_2) = \Delta W \quad (11)$$

Note that the change of potential energy Π is just the work done by conservative forces from r_1 to r_2 , namely

$$W = -(\Pi_2 - \Pi_1) = -\Delta \Pi \quad (12)$$

By (11), one can see

$$F = \nabla W = -\nabla \Pi \quad (13)$$

where ∇ is gradient operator.

Thus the component F_d of F in direction d is

$$F_d = F \cdot d = -d \cdot \nabla \Pi = -\frac{\partial \Pi}{\partial d} = \frac{\partial W}{\partial d} \quad (14)$$

Substituting (14) into (1) yields

$$\delta W = -\left(\frac{\partial \Pi}{\partial q_1} \delta q_1 + \frac{\partial \Pi}{\partial q_2} \delta q_2 + \dots + \frac{\partial \Pi}{\partial q_N} \delta q_N\right) = -\delta \Pi \quad (15)$$

If F can be decomposed into an internal force F_i and external forces F_e , then the work can also be set into internal work W_i and external work W_e . And then the total potential energy Π (TPE) is defined as

$$\Pi = U + V \quad (16)$$

where $U = -W_i$ is internal potential and $V = -W_e$ is external potential energy.

And

$$\delta \Pi = \left(\frac{\partial \Pi}{\partial q_1} \delta q_1 + \frac{\partial \Pi}{\partial q_2} \delta q_2 + \dots + \frac{\partial \Pi}{\partial q_N} \delta q_N\right) \quad (17)$$

is also defined the first variation of total potential energy Π .

By the minimum potential energy theorem (MPET), a conservative mechanical system is in a stable state of equilibrium if and only if the total potential energy is a relative minimum. By variational principle of stationary point, the MPET is equivalent to the statement: a system is in equilibrium if and only $\delta \Pi = 0$ for each and every virtual displacement, which is called mechanical variational theorem (MVT).

Comparing to (4), it is known that the MVT is actually a variational form in TPE. Hence we finish demonstrating that the VWM is a special form of VariM.

2.1.2 Complementary Virtual Work Method(CVWM)

Complementary virtual work is virtual work δW done by virtual force δF along a real displacement r .

$$\delta W = \delta F \cdot r \quad (18)$$

With respect to PVW, principle of complementary virtual work (PCVM) says that, the necessary and sufficient condition for a system to be in static equilibrium is that zero complementary virtual work is done for all virtual forces acting on the system, namely

$$\delta W = \sum_i \delta F_i \cdot r_i = 0 \quad (19)$$

where $\delta F_i, r_i$ are virtual forces and displacements respectively.

In a conservative system, total complementary potential energy (TCPE) is defined by

$$\Pi^* = U^* + V^* \quad (20)$$

where $U^* = -W_i$ is complementary internal potential and $V^* = -W_e$ is complementary external potential energy.

And by (18), the first variation of total complementary potential energy is

$$\delta \Pi^* = \sum_i \frac{\partial \Pi^*}{\partial F_i} \delta F_i \quad (21)$$

Like the MVT, it can derive that $\delta \Pi^* = 0$ is the sufficient and necessary condition for a conservative mechanical system to be in a stable state of equilibrium. Obviously, $\delta \Pi^* = 0$ is a variational form of the total complementary potential energy Π^* and is of course a special variational form of VariM..

3 Comments and Expectation

3.1 Conclusions About E_M and VariM

Deductions in section 2 show that the E_M, which is essentially PVW, is an instance of the VariM. The Total potential energy or Total complementary energy is the functional for the variation, which is demonstrated by (17) and (21). Therefore we say that, the E_M is a special form of the VariM while the inverse is false. This can be illustrated with Fig 2.

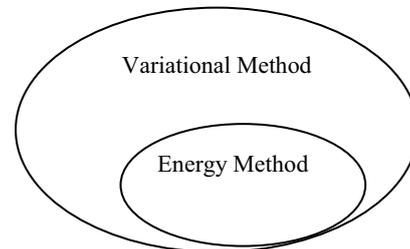


Fig 2. E_M is a subset of VariM

The relationship between the energy method and the variational method discloses a basic fact that foundation of mechanical computations is intrinsically related, and they are just “all roads lead to home”. The most importance of the intrinsic relation is that it is in some way to ensure the loyalty of the computational results to the computed objectives.

3.2 About Weighted Residual Method (WRM)

Fig 1 shows that the WRM is more general a method than the VariM and the E_M. In fact, anyone who has gotten an excellent background of computational mechanics knows that, the WRM can deduce more concrete computational methods than the VariM. For example, the WRM has deduced collocation method, sub-domain method, least square method, momentum method and even the Galerkin method while the VariM has up to now deduced the Ritz method. Nevertheless, the WRM still intersects with the VariM owing to the fact that the Galerkin method of the WRM can derive the same form of governing equation as that of the Ritz method of the VariM if the variation of a functional exists. Namely, the WRM intersects with the VariM via the Galerkin-Ritz method. Fig 3 demonstrates the relationship.

Another fact existing in computational mechanics is that the VariM has strong mathematical foundations and the E_M can control the lower and upper bound of

computations in an appreciate range while people up to now have no way to control the bound of the residual. This may be the reason that people prefer using the VariM and the E_M to the WRM, particularly in computational solid mechanics.

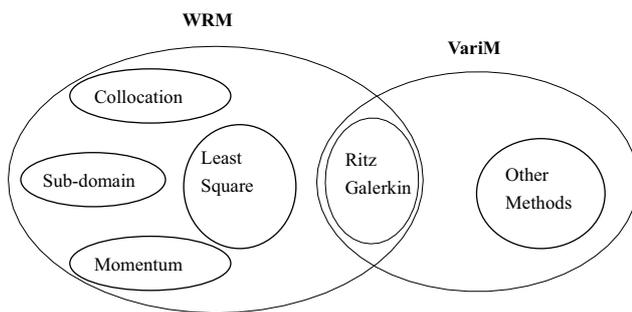


Fig 3. Ritz-Galerkin mthod bridegs WRM and VariM

3.3 Expectation

Up to now it has been known that the energy method, the variational method and the weighted residual method are three fundamental methods in computational mechanics and the energy method is essential a special form of the variational method. It is also known that the weighted residual method is more general and incorporates with the variational method in some way, e.g., the Ritz-Galerkin method. This situation leaves a lot of imagination and work for us to do. On the one hand, the energy method is a subset of the variational method, and on the other hand the variational method has gotten an intersection with the weighted residual method. Now we expect to know if this intersection can be enlarged under some universal principle. If it can, what is the universal principle? If it cannot, what are the barriers? The answers to these questions need more and deeper researches in foundation of computational mechanics. It may concerns more intrinsic essence of mechanics and might bring more ingenious results for mechanics, which are our hopes.

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